# FORCE FIELD STUDY OF SOME XY2 TYPE MOLECULES

BY K. RAMASWAMY AND R. SRINIVASAN

Department of Physics, Annamalai University, Annamalainagar\*

(Received December 9, 1975; final version received May 4, 1976)

The most probable force fields of some of the  $XY_2$  type molecules like  $SO_2$ ,  $NO_2$ ,  $SiF_2$  and  $CF_2$  have been fixed using the method of Redington and using the observed molecular parameters such as the Coriolis coupling constant, mean amplitudes of vibration, rotation distortion constants and isotopic shifts.

### 1. Introduction

It is well known that the unambiguous determination of force fields is possible only with the help of some additional data such as the Coriolis coupling constants [1], mean amplitudes of vibration [2], rotation distortion constants [3], and isotopic shifts [4]. Attempts have been made to obtain mathematically feasible and physically meaningful set of torce constants using several methods [1-6]. A critical survey of all these methods in uniquely fixing the force field was made earlier for sulfur trioxide [2]. The present paper is an extension of the above method to fix the unique force fields of some of the  $XY_2$  bent type molecules namely  $SO_2$ ,  $NO_2$ ,  $SiF_2$  and  $CF_2$ .

## 2. Spectral data

The necessary vibrational frequencies and molecular parameters used in the present calculations are summarised in Table I along with the references [7–19] from which they are taken. The symmetry coordinates used here are essentially the same as those reported by Cleveland and Meister [20].

<sup>\*</sup> Address: Department of Physics, Annamalai University, Annamalainagar 608101 Tamilnadu, S. India.

The observed vibrational frequencies (cm<sup>-1</sup>), bond parameters, Coriolis coupling constants, rotation distortion constants (MHz) and mean amplitudes of vibration (Å) of some XY<sub>2</sub> type molecules

Molecule	Vibrational frequencies			Bond parameters		Coriolis coupling constant	Rotation distortion contants in MHz.		
	$v_1(A_1)$	$v_2(A_1)$	ν <sub>3</sub> (B <sub>2</sub> )	$R_{x-y}$	XŶX	ζ13	$ au_{XXXX}$	$ au_{YYYY}$	$ au_{XXYY}$
<sup>32</sup> S <sup>16</sup> O <sub>2</sub>	1156 [7]	522	1366	1.435 [11]	119° <b>2</b> 1′	0.3066 [14]	-9.9216 [17]	-0.0399	0.4314
<sup>34</sup> S <sup>16</sup> O <sub>2</sub>	1147 [7] 1357.8 [8]	518 756.8	1349 1665.5	1.197 [12]	134°15′	<u></u>	<del>-299.4</del> [18]	-0.0414	1.8430
<sup>15</sup> N <sup>16</sup> O <sub>2</sub> <sup>28</sup> Si <sup>19</sup> F <sub>2</sub>	1342.5 [8] 843 [9]	747.1 343	1628 855	1.5946 [11]	100°53′	0.4285 [15]	-1.9807 [19]	-0.0715	0.2613
$^{12}C  ^{19}F_2$ $^{13}C  ^{19}F_2$	1222 [10] 1191 [10]	668 668	1102 1073	1.3035 [13]	104°47′		-11.2778 [13]	-0.0691	

Note: The electron diffraction values of mean amplitudes of vibration of both bonded and nonbonded atoms are  $\sigma_{x-y} = 0.041$  and  $\sigma_{y...y} = 0.053$ , Ref. [16].

### 3. Results and discussion

The Redington method [5] was successfully applied to the SO<sub>2</sub>, NO<sub>2</sub> and SiF<sub>2</sub>. "F<sub>stp</sub>" solutions of 47°15′, 53°15′ and 27°45′ were obtained by this method. But for CF<sub>2</sub>, Redington's method could not yield reliable force field. The observed Coriolis coupling constants of SO<sub>2</sub> and SiF<sub>2</sub> were used as additional data and angles of 48°30′ and 28°30′ were determined which reproduced the experimentally observed molecular constants.

Electron diffraction values of the mean amplitudes of vibration were available only for the  $SO_2$  molecule. The use of mean amplitude of vibration of non bonded atoms yields a rather low value for  $\phi$ , and as a result does not reproduce any of the molecular constants. The mean amplitude of vibration of bonded atoms gives only an imaginary solution.

The rotation distortion constants like  $\tau_{XXXX}$ ,  $\tau_{YYYY}$  and  $\tau_{XXYY}$  give comparable force fields for NO<sub>2</sub> and SiF<sub>2</sub>. In the case of SO<sub>2</sub> the  $\tau_{XXXX}$  and  $\tau_{YYYY}$  parameters alone yielded  $\phi$ values which coincide well with the values determined by other methods. But the  $\tau_{XXYY}$  experimental value is higher than those calculated for various angle parameters varying from 0° to 180° and hence it could not be used to fix the force field. For CF<sub>2</sub>, the  $\tau_{XXXX}$  and  $\tau_{YYYY}$  parameters yielded a very reliable set of force constants and other molecular constants.

The isotopic shifts used for the NO<sub>2</sub> and CF<sub>2</sub> furnished force fields comparable with the force fields obtained through the other methods. But for SO<sub>2</sub> since the angle obtained was very low, the molecular constants determined for this parameter were considerably lower than the others.

The various  $\phi$  values obtained through different methods are given in Table II.

The  $\phi$  values obtained through different methods

Method	Parameter $\phi$					
Weinod	SO <sub>2</sub>	NO <sub>2</sub>	SiF <sub>2</sub>	CF <sub>2</sub>		
Redington	47°15′	53°15′	27°45′	15°46′		
Using 5	48°30′		28°30′			
Using $\sigma_{YY}$	42°10′					
Using $\tau_{XXXX}$	51°	53°	29°	27°32′		
Using $ au_{YYYY}$	47°15′	53°15′	29°30′	25°51′		
Using $ au_{XXYY}$	_	53°36′	27°45′			
Using isotopic shifts	43°41′	51°16′	_	25°15′		

Most probable  $\phi$ : The most probable  $\phi$  can be stated as the average of all the  $\phi$ 's obtained by different approaches which reproduce all molecular constants within reasonable accuracy. The various molecular constants obtained for the most probable  $\phi$  of  $SO_2$ ,  $NO_2$ ,  $SiF_2$  and  $CF_2$  are given in Table III.

TABLE III Force constants (mdyn/Å), Coriolis coupling constants, mean amplitudes of vibration (Å) and rotation distortion constants (MHz) obtained for the most probable  $\phi$  of some  $XY_2$  type molecules

Molecular	Molecule							
constants	SO <sub>2</sub>	NO <sub>2</sub>	$SiF_2$	NF <sub>2</sub>				
Most probable ∮	47°40′	52°53′	28°30'	26°13′				
$F_{11}$	10.0793 <sup>a</sup> (10.03) [21] (10.05) [22]	13.1265 (12.963) [8] (13 185) [18]	5.1636 (5.329) [15] (5.34) [25]	7.656 (7.756) [13] (7.450) [26]				
F <sub>22</sub>	0.8086 (0.7933) [21] (0.7930) [22]	1.1181 (1.125) [8] (1.109) [18]	0.4356 (0.44) [15] (0.439) [25]	1.383 (1.373) [13] (1.400) [26]				
F <sub>12</sub>	0.2513 (0.267) [21] (0.280) [22]	0.6281 (0.055) [8] (0.0679) [18]	0.1634 (0.174) [15] (0.195) [25]	0.7413 (0.7713) [13] (0.6800) [26]				
ζ <sub>13</sub>	0.3224	0.4755	0.4285	0.7182				
$\sigma_{X-Y}$	0.0351	0.0381	0.0426	0.0450				
$\sigma_{\mathbf{Y}\mathbf{Y}}$	0.0507	0.0431	0.0745	0.0505				
$ au_{XXXX}$	-9.3029	-299.035	-1.9492	-10.9343				
$ au_{YYYY}$	-0.0393	-0.0417	-0.0729	-0.0689				
$ au_{XXYY}$	0.3979	1.8488	0.2637	_ *				

<sup>&</sup>lt;sup>a</sup> This number of significant figures is retained to secure internal consistency in the calculations.  $mdyn/Å = 10^2 N/m$ . The values in parenthesis are taken from their respective references.

For the SO<sub>2</sub> molecule, as in Redington's method, we observed the Coriolis coupling constant and  $\tau_{XXXX}$  and  $\tau_{YYYY}$  parameters furnished comparable force fields. An average value of the three was taken as the most probable  $\phi$ .

The force constants calculated for the most probable  $\phi$  of 47°40′ compare extremely well with those reported by Kivelson [21], Polo and Wilson [22] and Morino et al. [14]. The mean amplitudes of vibration of both bonded and non-bonded atoms determined for the most probable  $\phi$  coincide well with those spectroscopically calculated by Clark and Beagley [16], even though their experimental values are higher than these values. But the most probable  $\phi$  could not reproduce the Coriolis coupling constant accurately. This may probably be due to the fact that  $\xi$  is a highly sensitive parameter in fixing the force fields. The rotational distortion constants are in good agreement with the reported values.

For the NO<sub>2</sub>, the most probable  $\phi$  was determined as the average of the angle parameters obtained through Redington's method, using observed rotation distortion constants and isotopic shifts. The force field furnished by this angle of 52°53′ can be assumed to be unique as it reproduces the diagonal force constants reported by Arkawa et al. [8], Bird et al. [18], the observed rotation distortion constants and Coriolis coupling constants calculated by Cyvin [23], namely  $\xi_{13}^2 = 0.249$  and Oka and Morino [24] namely  $\xi_{13}^2 = 0.241$ . The off diagonal force constants differ considerably from the reported values.

In the case of SiF<sub>2</sub>, the force constants reported by Khanna et al. [15] and Shoji and others [25] were accurately reproduced by the most probable force field in addition to the observed Coriolis coupling constants and rotation distortion constants.

The Redington method could not yield a reliable force field for the CF<sub>2</sub> molecule. The experimental isotopic shifts and rotation distortion parameters were utilized in uniquely fixing the force field. This most probable force field yields a good set of force constants in agreement with the values of Milligan et al. [26] and Kirchoff et al. [13].

One of the authors (R. S.) is grateful to the Council of Scientific and Industrial Research, Government of India, New Delhi, for financial assistance in the form of a Junior Research Fellowship.

#### REFERENCES

- [1] T. R. Ananthakrishnan, G. Aruldhas, J. Mol. Struct. 13, 1636 (1972).
- [2] K. Ramaswamy, R. Srinivasan, Acta Phys. Pol. A49, 763 (1976).
- [3] L. H. Jones, R. R. Rayan, L. B. Asprey, J. Chem. Phys. 49, 581 (1968).
- [4] B. Jordonov, B. Nickolova, J. Mol. Struct. 13, 21 (1972); 15, 7 (1973); 15, 19 (1973).
- [5] R. L. Redington, A. L. Khidir Aljibury, J. Mol. Spectrosc. 37, 494 (1971).
- [6] E. B. Wilson Jr., J. C. Decius, P. C. Cross, Molecular Vibrations, McGraw Hill, New York 1955.
- [7] M. Allawena, R. Rysnik, D. White, V. Calder, D. E. Mann, J. Chem. Phys. 50, 3399 (1969).
- [8] E. T. Arkawa, A. H. Nielson, J. Mol. Spectrosc. 2, 413 (1958).
- [9] D. E. Milligan, M. E. Jacox, J. Chem. Phys. 49, 4269 (1968).
- [10] D. E. Milligan, D. E. Mann, M. E. Jacox, R. A. Mitch, J. Chem. Phys. 41, 1199 (1964).
- [11] W. H. Kirchoff, J. Mol. Spectrosc. 41, 333 (1972).
- [12] G. R. Bird, R. Kato, J. Rolfe, J. Chem. Phys. 25, 1040 (1968); 47, 1901 (1967).
- [13] W. H. Kirchoff, D. R. Lide Jr., F. X. Powell, J. Mol. Spectrosc. 47, 491 (1973).

- [14] Y. Morino, Y. Kikuchi, S. Saito, E. Hirota, J. Mol. Spectrosc. 13, 95 (1964).
- [15] V. M. Khanna, R. Hauge, R. F. Curl Jr., J. L. Margrave, J. Chem. Phys. 47, 5031 (1967).
- [16] A. H. Clark, B. Beagley, Trans. Faraday Soc. 67, part 8 (1971).
- [17] H. A. Gebbie, W. B. Stone, E. K. Gora, S. A. Clough, F. X. Kneizys, J. Mol. Spectrosc. 19, 7 (1966).
- [18] G. R. Bird, J. C. Baird, A. W. Jache, J. A. Hodgeson, R. F. Curl Jr., A. C. Kunkee, J. W. Bransford, J-Anderson Rastrup, J. Rosenthal, J. Chem. Phys. 40, 3378 (1966).
- [19] V. M. Rao, R. R. Curl Jr., J. Chem. Phys. 45, 2032 (1966).
- [20] A. G. Meister, F. F. Cleveland, Am. J. Phys. 41, 13 (1946).
- [21] D. Kivelson, J. Chem. Phys. 22, 904 (1954).
- [22] S. R. Polo, M. K. Wilson, J. Chem. Phys. 22, 900 (1954).
- [23] S. J. Cyvin, J. Mol. Spectrosc. 11, 195 (1963).
- [24] T. Oka, Y. Morino, J. Mol. Spectrosc. 8. 9 (1962).
- [25] H. Shoji, T. Tawaka, E. Hirota, J. Mol. Spectrosc. 47, 268 (1973).
- [26] D. E. Milligan, M. E. Jacox, J. Chem. Phys. 48, 2265 (1968).